Gamma-ray burst spectral evolution in the internal shock model

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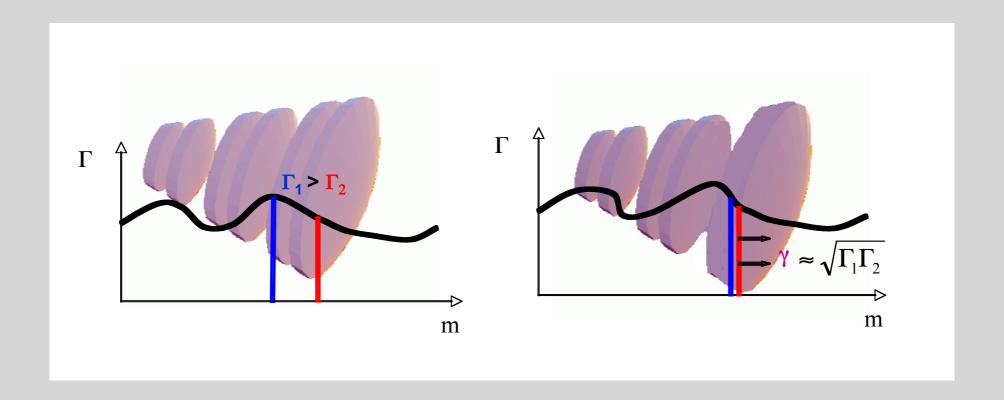
in collaboration with:

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Prompt high energy emission in the framework of internal shocks

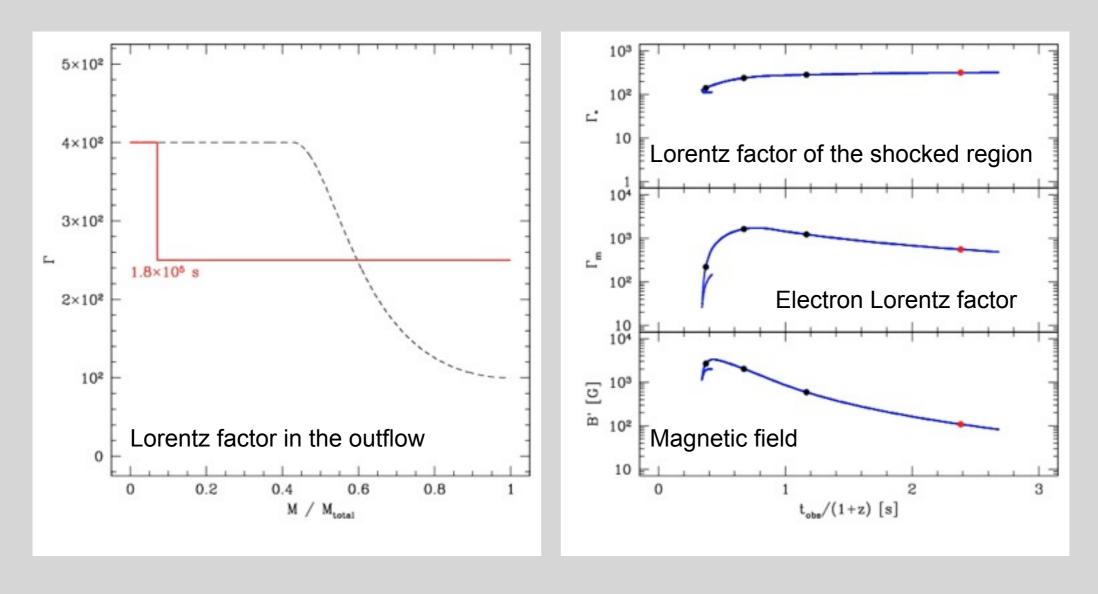


Modeling:

- 1. dynamics of internal shocks
- 2. radiative processes in the shocked medium
- 3. observed spectra and time profiles

Dynamics of the internal shocks

Physical conditions in the shocked medium: Lorentz factor $\Gamma*$, comoving density ρ^* , comoving specific energy density ϵ^*



Dissipated energy is distributed between protons, electrons (fraction ε_{e}) and magnetic field (fraction ε_{e})

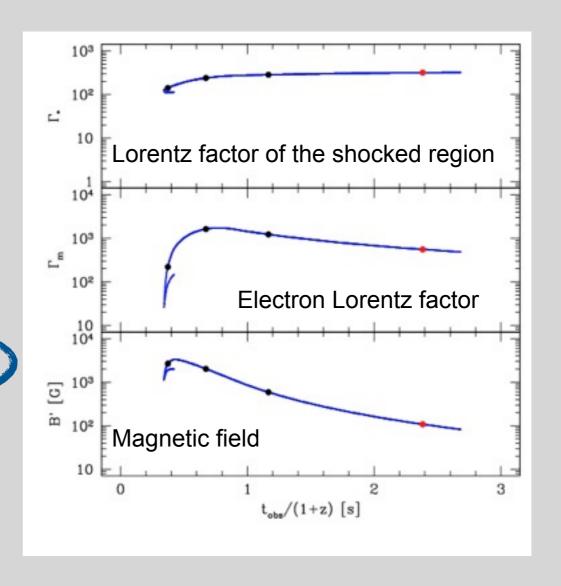
Dynamics of the internal shocks

Physical conditions in the shocked medium: Lorentz factor $\Gamma*$, comoving density ρ^* , comoving specific energy density ϵ^*

Relativistic electron density:

$$n'(\Gamma_e, t'=0) \propto \Gamma_e^{-p} \qquad \Gamma_e \geq \Gamma_m$$

 ξ < | of all electrons is accelerated



Dissipated energy is distributed between protons, electrons (fraction ε_{e}) and magnetic field (fraction ε_{e})

Radiative processes

Assumption: instantaneous shock acceleration

Adiabatic cooling timescale: $t'ex = R / \Gamma^* c$ (comoving frame)

Radiative timescale: **t**`rad

t'rad << t'ex high radiative efficiency

Electron and photon distributions evolve strongly with time!

The present version of the code follows the time evolution of the electron density and the photon density including the following processes:

- adiabatic cooling (spherical expansion)
- synchrotron
- inverse Compton
- synchrotron self-absorption
- $\gamma\gamma$ annihilation

Not included:

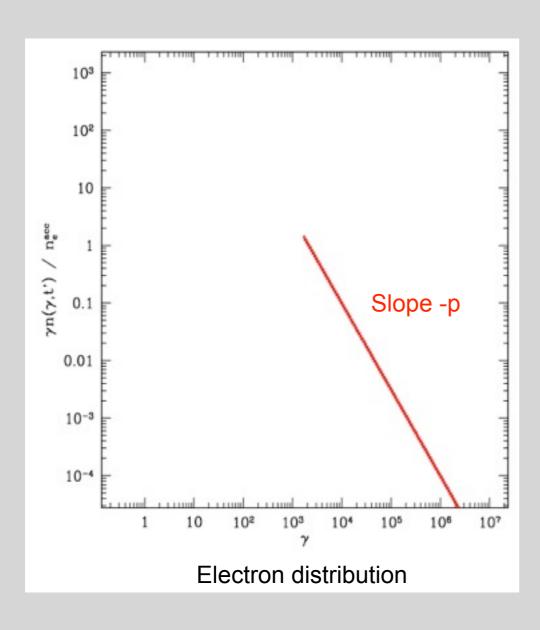
- * emission from secondary leptons
- * IC in optically thick regime (Comptonisation)

ELECTRONS:
$$\frac{\partial n'}{\partial t'}(\Gamma'_{e},t') = -\frac{\partial}{\partial \Gamma'_{e}} \left[\left(\frac{d\Gamma'_{e}}{dt'} \Big|_{syn+ic} + \frac{d\Gamma'_{e}}{dt'} \Big|_{ad} \right) n'(\Gamma'_{e},t') \right]$$

$$\frac{\partial n'_{\mathbf{v}}}{\partial t'} = \int n'(\Gamma'_{e}, t') P_{syn+ic}(\Gamma'_{e}) d\Gamma'_{e} - cn'_{\mathbf{v}} \int n'(\Gamma'_{e}, t') \sigma_{abs}(\Gamma'_{e}, \mathbf{v}) d\Gamma'_{e} - cn'_{\mathbf{v}} \int_{\mathbf{v}' > \frac{(m_{e}c^{2})^{2}}{h^{2}\mathbf{v}}} n'_{\mathbf{v}'}(t') \sigma_{\gamma\gamma}(\mathbf{v}, \mathbf{v}') d\mathbf{v}'$$

Radiative processes

Radiation: the time evolution of electrons and photons in the comoving frame is solved (time-dependent radiative code)



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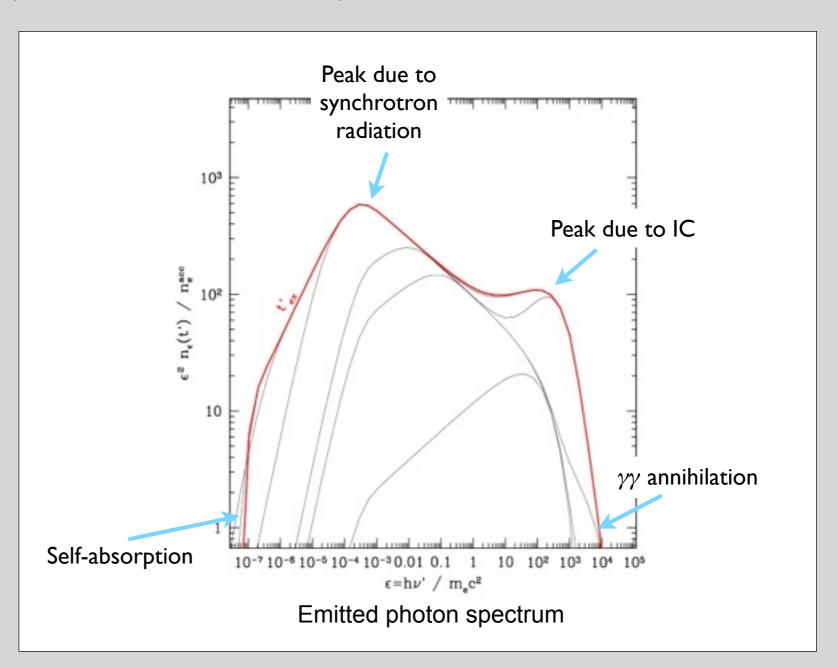
Comptonization parameter Y = Lic / Lsyn

IC dominant:

Iow frequency synchrotron peak Thomson regime

Synchrotron dominant:

high frequency synchrotron peak Klein-Nishina regime



This calculation is done at all times along the propagation of each shock wave All the contributions are added together to produce a synthetic gamma-ray burst (spectrum+lightcurve)

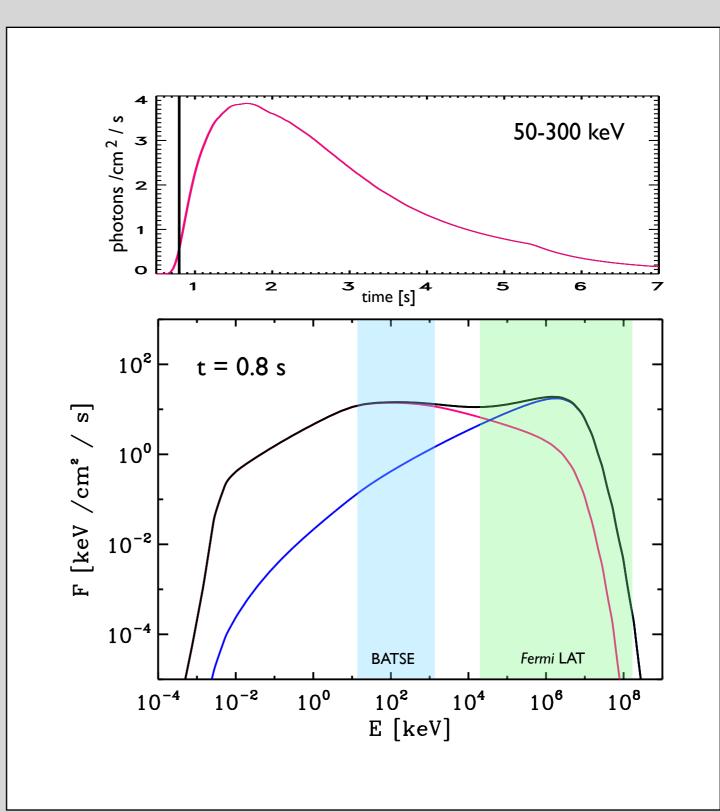
Observed spectra and time profiles

The observed spectra and the light curves are computed from the comoving emission by integration over equal-arrival time surfaces.

relativistic effects
(Doppler factor)
geometry (curvature of the
emitting surface)
cosmological effect (redshifts)

Instantaneous observed spectrum:

synchrotron inverse Compton total



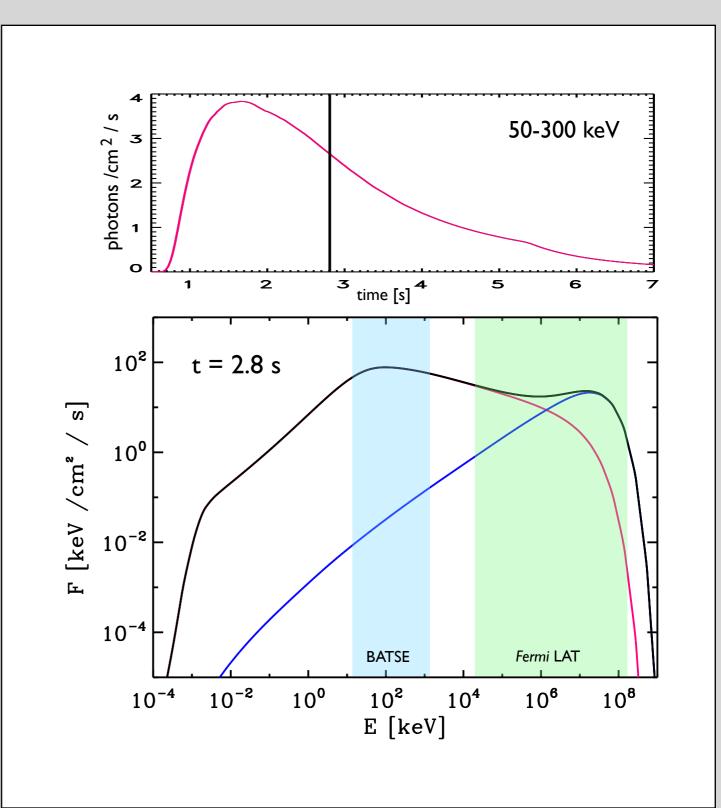
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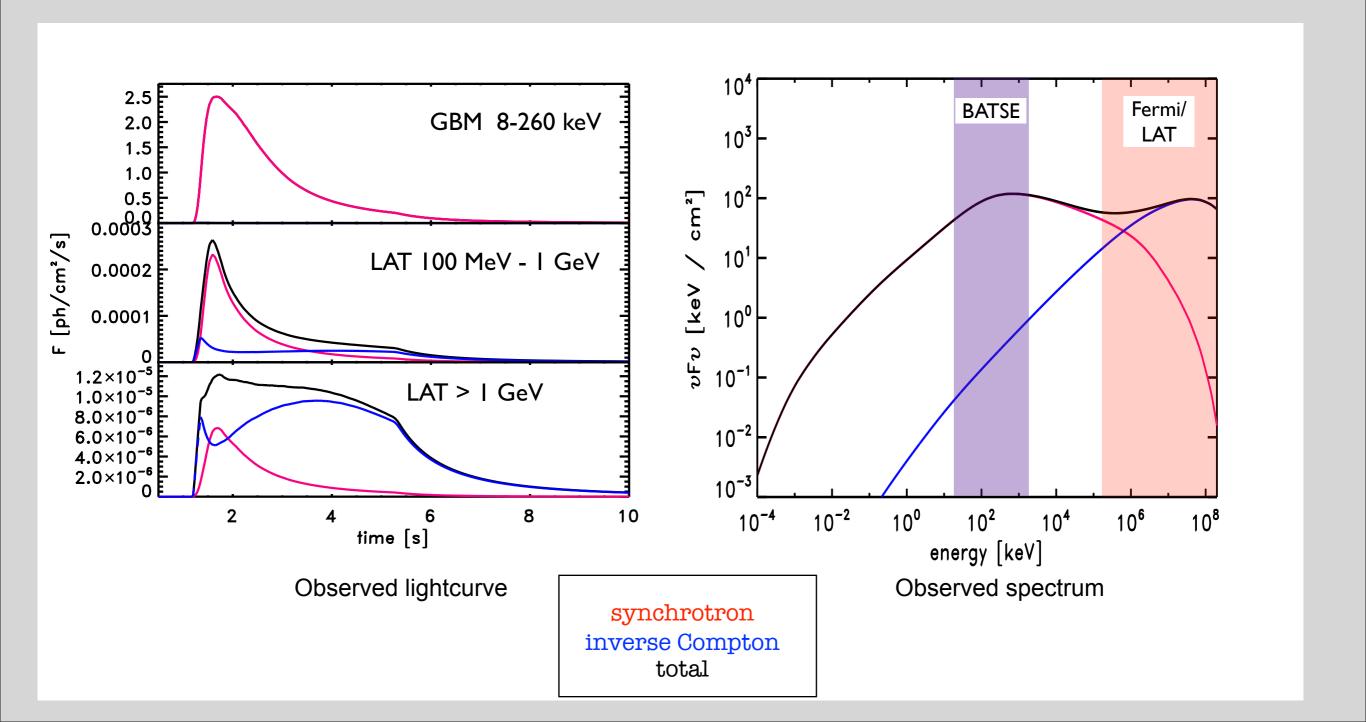


Dominant radiative process in sub-MeV range?

SYNCHROTRON CASE (B)

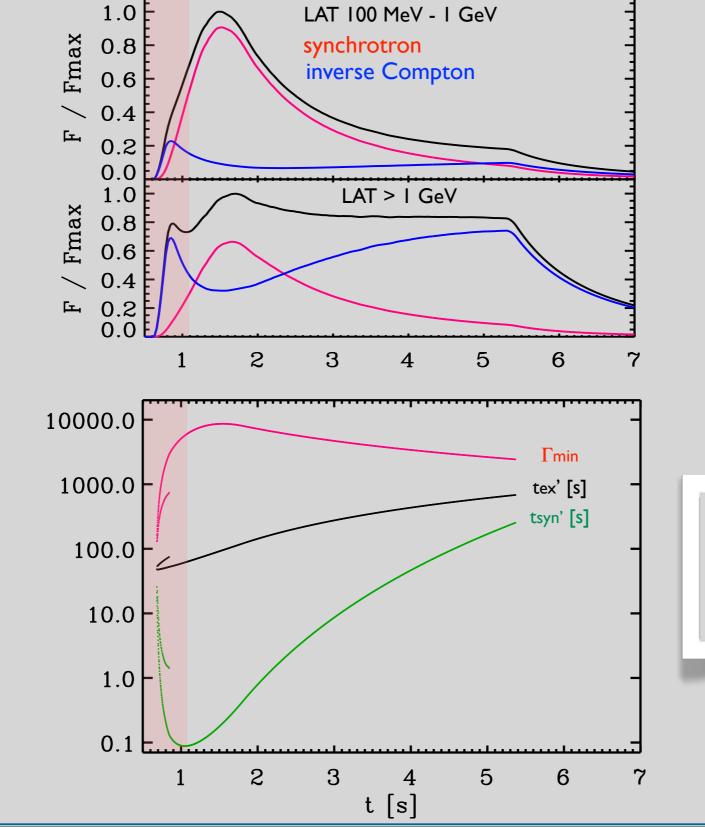
low magnetic field

dE/dt = 5 x 10 ⁵³ erg s,
$$\varepsilon_{\rm B}$$
 = 0.003, $\varepsilon_{\rm e}$ = 1/3, ζ = 0.003, p = 2.5, z=1



Temporal profiles: > 100 MeV bands

Model: in LAT (>100 MeV) energy bands both components present, synchrotron + IC



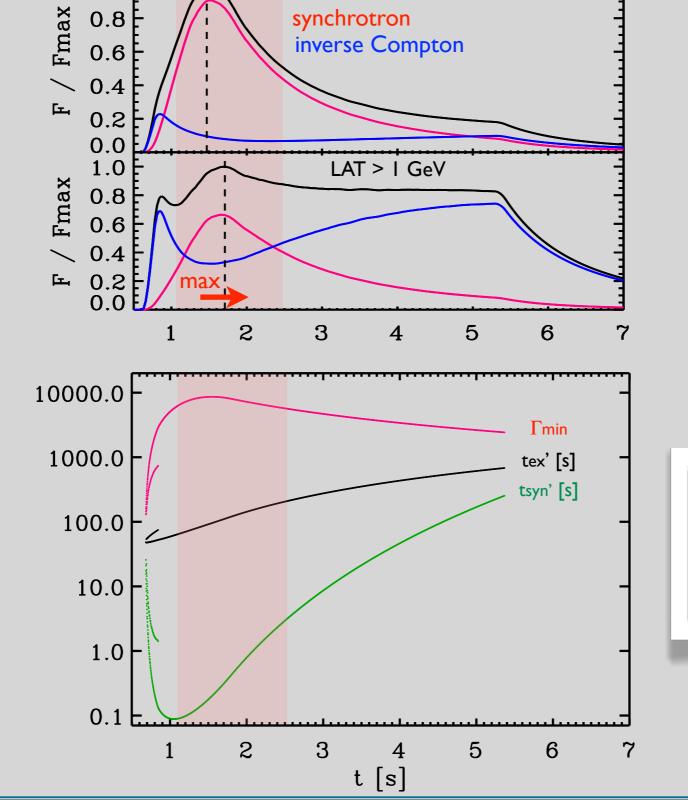
 $\begin{array}{c} \text{weak shock} \\ \epsilon^* \text{ low} \\ \text{moderate } \Gamma_\text{m} \Rightarrow \text{large tsyn'} \\ \text{R small} \Rightarrow \text{tex'} \cong \text{R/}\Gamma^*\text{c small} \\ \text{tsyn'} \leq \text{tex'} \Rightarrow \text{large efficiency of IC} \end{array}$

Temporal profiles: > 100 MeV bands

1.0

Model: in LAT (>100 MeV) energy bands both components present, synchrotron + IC

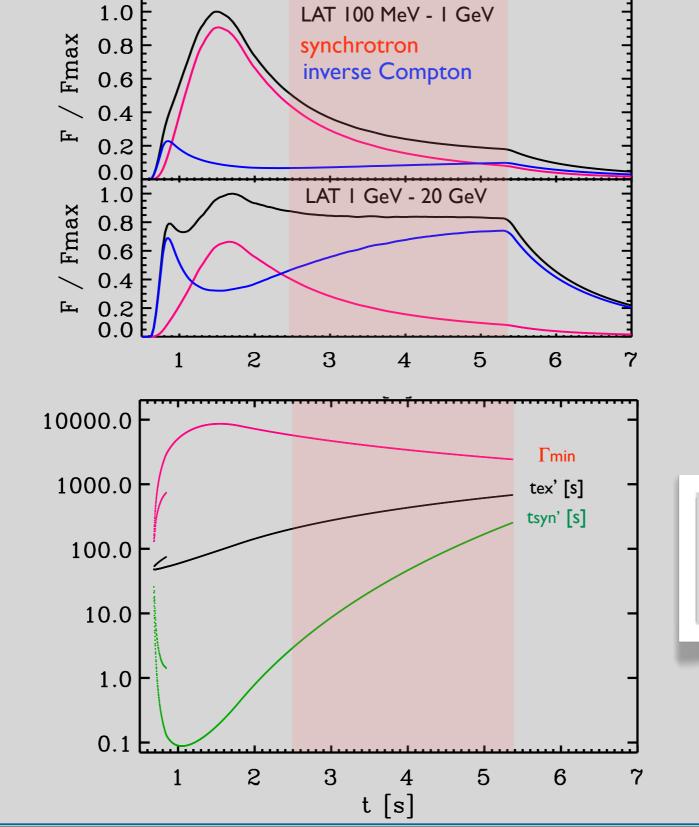
LAT 100 MeV - I GeV



shock becomes stronger Γ m increases \Rightarrow tsyn' decreases R, tex' increase tsyn' << tex' \Rightarrow low efficiency of IC dominant synchrotron component

Temporal profiles: > 100 MeV bands

Model: in LAT (>100 MeV) energy bands both components present, synchrotron + IC



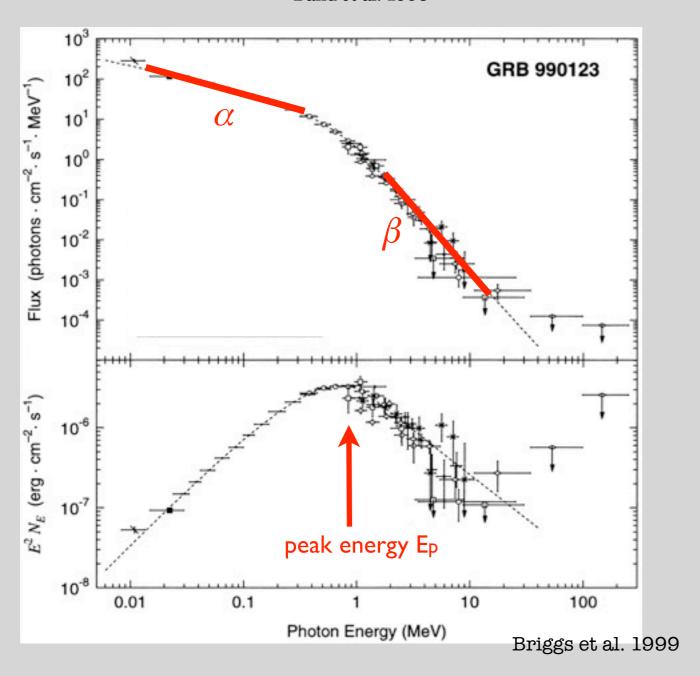
tail of the pulse:

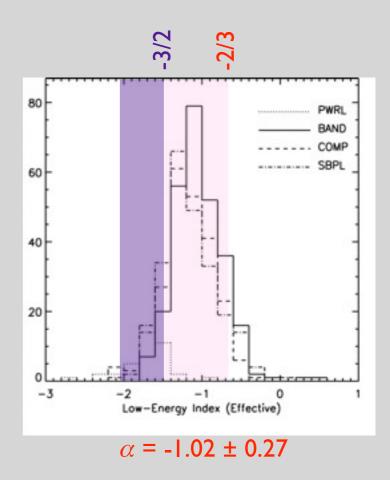
B decreases ⇒ tsyn' increases

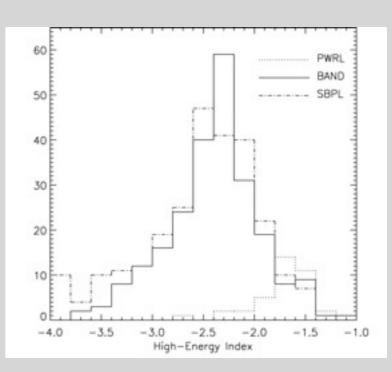
tsyn' ≤ tex' ⇒ increased efficiency of IC

IC component dominant in GeV

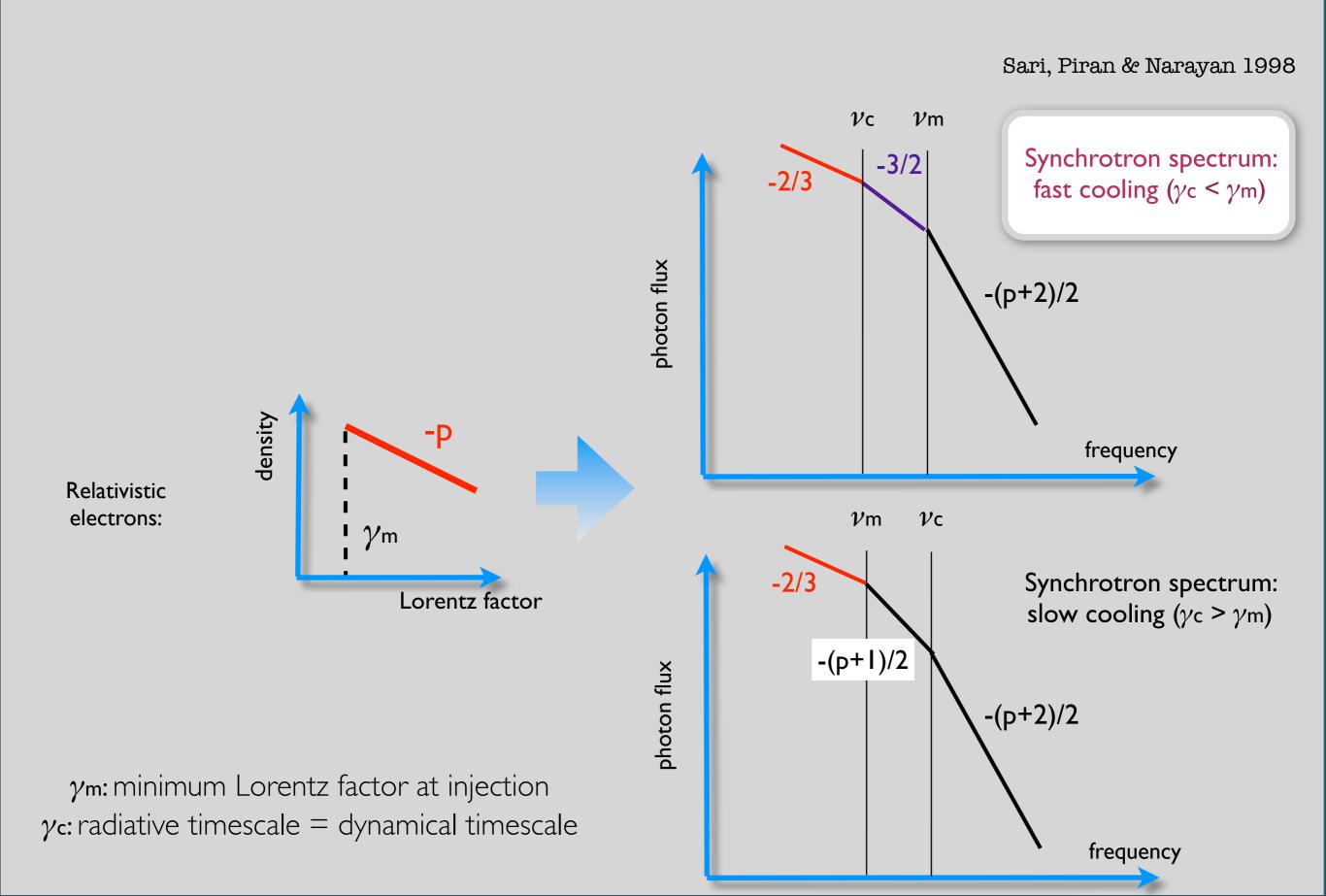
4-parameters "Band spectrum" E_P, α, β and normalization Band et al. 1993

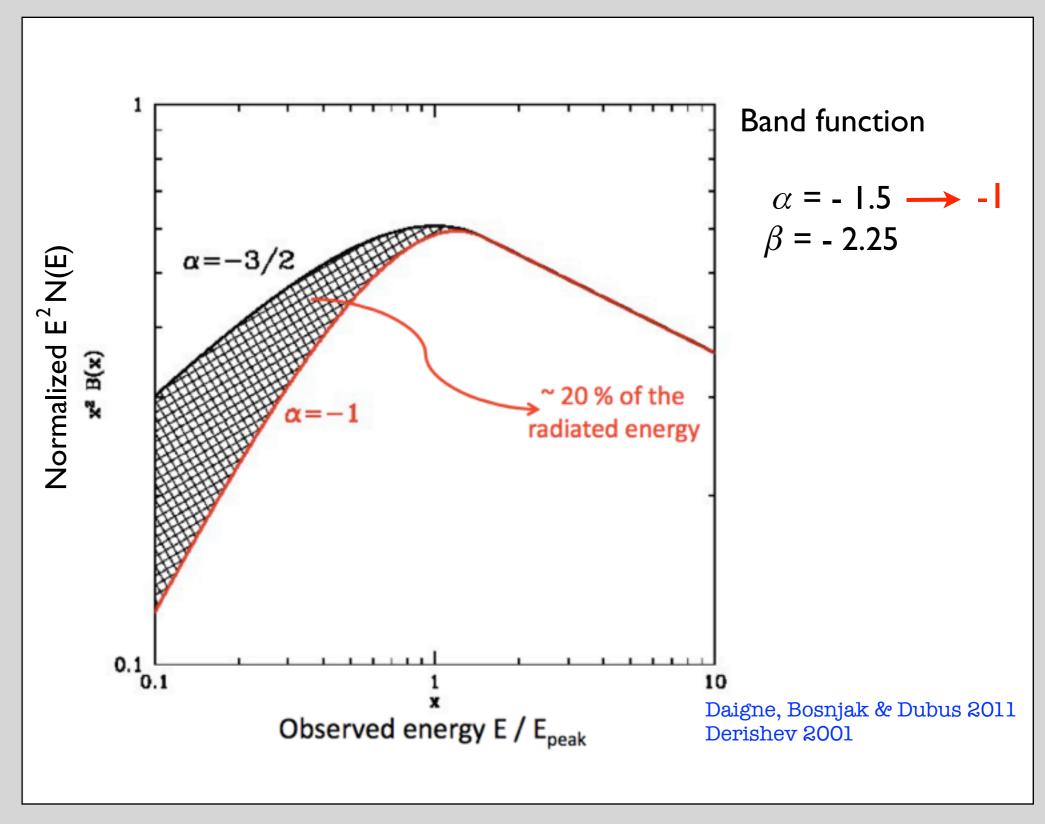




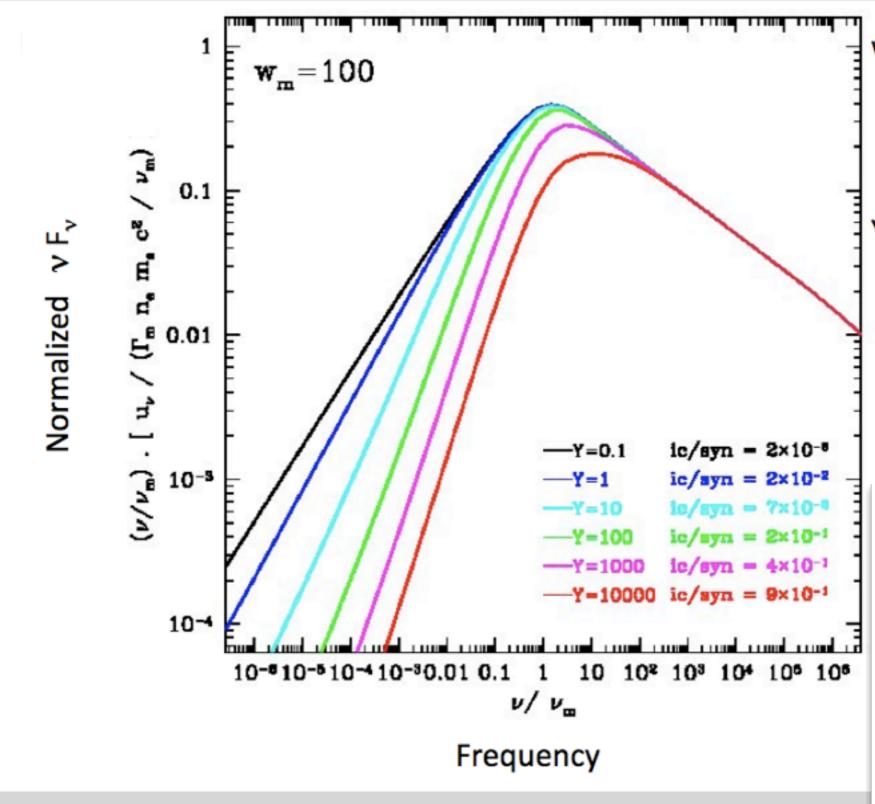


 β = -2.35 ± 0.27





Inverse Compton scatterings in Klein-Nishina regime have an impact on the synchrotron slope



w_m: importance of KN

$$w_{\rm m} = \Gamma_{\rm m} \frac{h \nu_{\rm m}'}{m_{\rm e} c^2}$$

Y: importance of IC vs syn

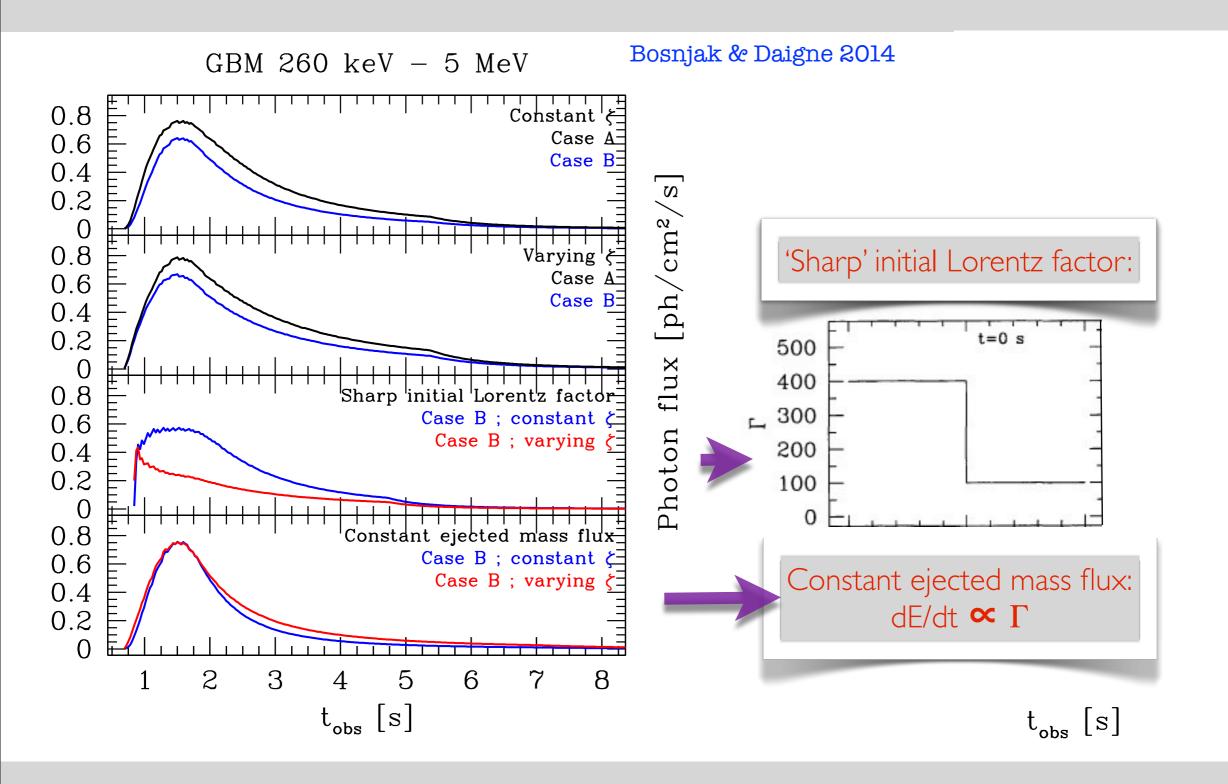
$$Y = rac{4}{3} au_{
m T} \Gamma_{
m m} \Gamma_{
m c} \simeq rac{\epsilon_{
m e}}{\epsilon_{
m B}}$$

Thomson regime: the electron cooling rate due to IC scatterings remains proportional to γ^2 as for the synchrotron power

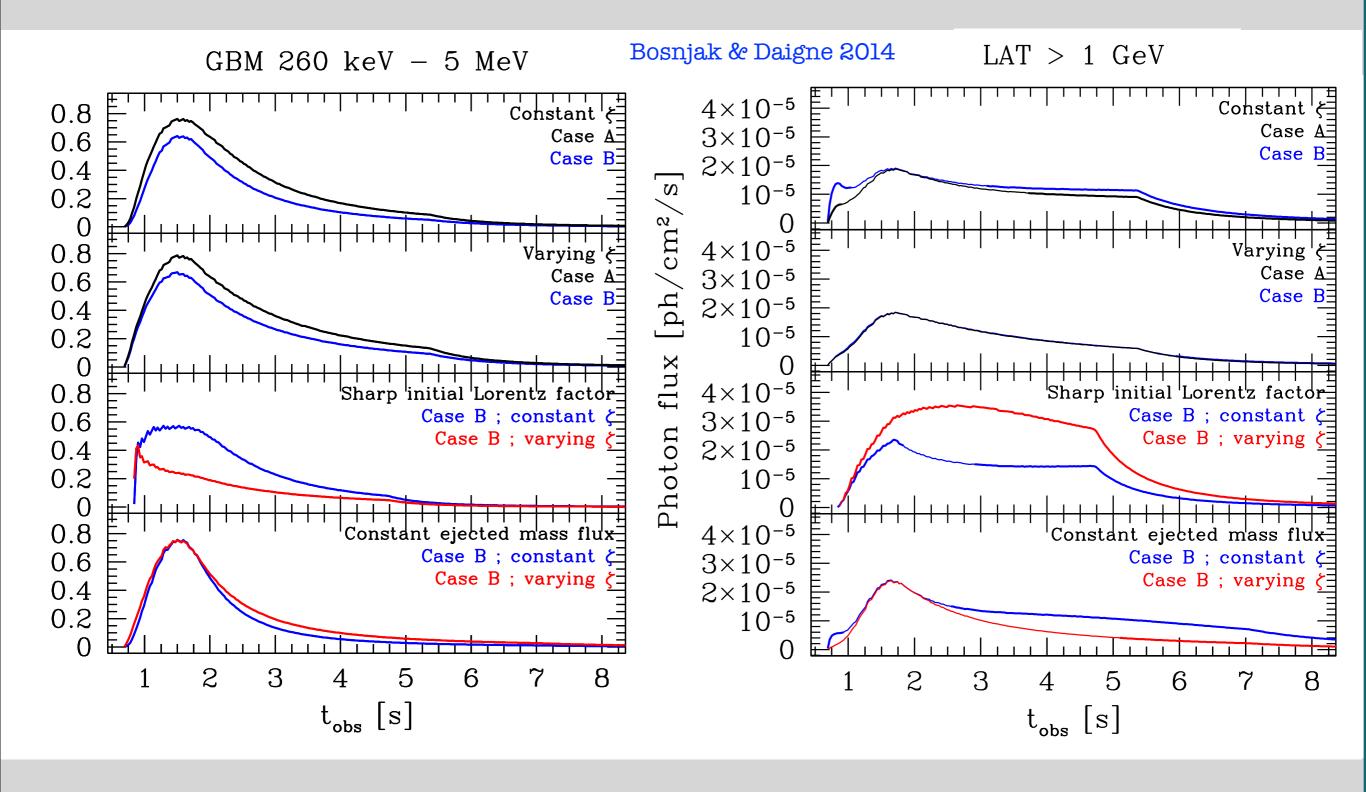
KN regime: the electron cooling rate due to IC depends on γ

Exact calculation with synchrotron + IC only (no adiabatic cooling, synchrotron self-absorption, $\gamma\gamma$ annihilation)

High energy emission: light curves



High energy emission: light curves



Summary

We developed modeling tools to compute the GRB prompt emission from internal shocks in a time-dependent way in different spectral bands, including the high-energy gamma rays

The exploration of the parameter space shows that we can expect two classes of broad-band spectra, which correspond to different physical conditions in the shocked region: the "synchrotron case" (where the dominant process in Fermi-GBM range is synchrotron radiation) and the "inverse Compton case" (where the synchrotron component peaks at low energy and the dominant process in the GBM range is inverse Compton)

Fermi GRB observations favor the "synchrotron case", with inverse Compton scatterings occurring in Klein-Nishina regime. This scenario qualitatively reproduces the observed spectral evolution (HIC, HFC). We constrain the parameters of the model (p, εB , ζ) in order to have a quantitative agreement

Further developments: currently incorporating a more realistic scenario for the physical conditions in the shocked plasma and making predictions for the CTA observatory