

Hypervelocity Stars

Tidal disruption of binaries by a massive BH

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Collaborators

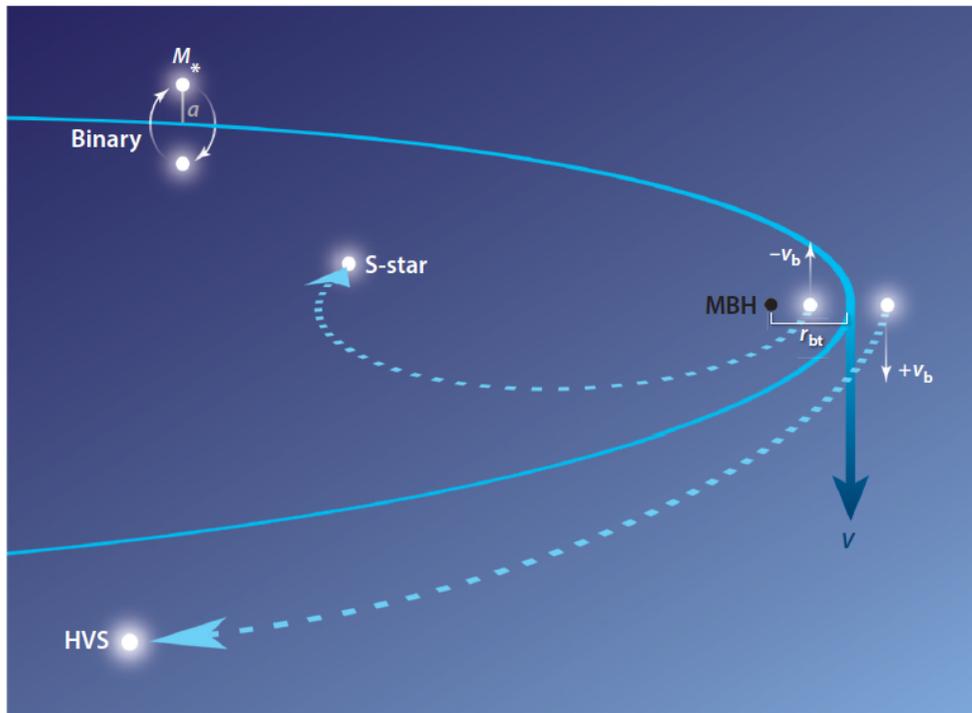
Harry Brown (LJMU) , Elena Rossi (Leiden), Re'em Sari (Hebrew)

Hypervelocity Stars (HVSs)

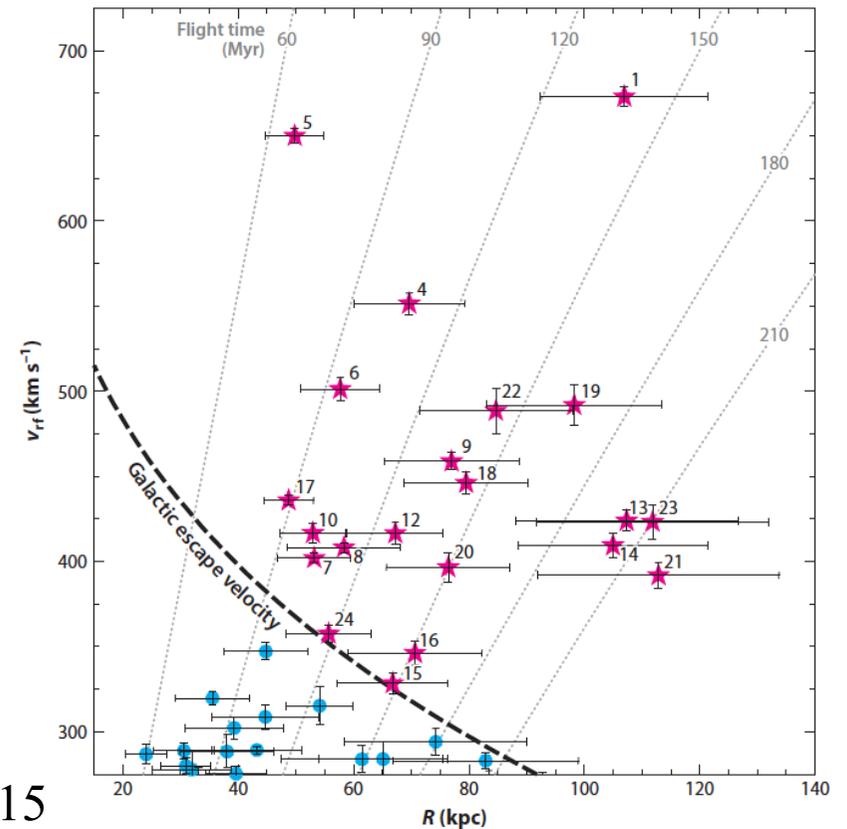
ejected by the Galaxy's central massive BH at speeds that exceed Galactic escape velocity

Prediction by Hills 1988: Tidal disruption of binaries by the massive BH

Discovery by Brown+ 2005: B-type stars with such high velocities in the halo



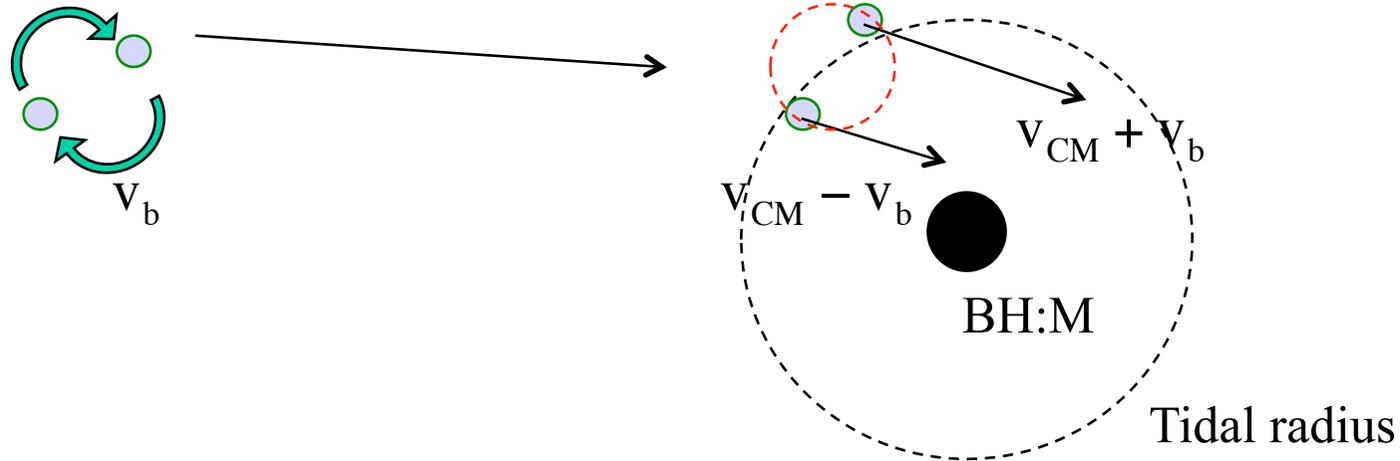
Brown 2015



Implications of HVSs

- Important tool to understand the nature and environment of massive BH/BHs
 - Single BH (isotropic ejection) vs BH binary (BH orbital plane)
 - Stellar populations (mass function, binary fraction, ages)
- Tracers of the Galactic Potential (stars & dark matter)
 - Shape and orientation of dark matter halo (Gnedin 2005)
 - Triaxiality of dark halo (Yu&Madau 2007)
- Links to TDEs and massive BH Growth
 - HVS ejection rate: $10^{-5} - 10^{-4}$ per year
 - Every HVS leaves its companion behind (S-stars).
 - Central BH may have grown owing to binary disruption and subsequent TDEs
 - a factor of 2-4 in past 5-10Gyr (Bromley+ 2012).

Ejection velocity estimate



Binary rotation velocity: $v_b \sim \sqrt{Gm/a}$.

Binary CM centre moves at r_t with $v_{CM} \sim \sqrt{GM/r_t} \sim v_b (M/m)^{1/3}$.

Additional kinetic energy $\Delta\epsilon \sim v_b v_{CM}$

$$v \sim \sqrt{\Delta\epsilon} \sim (M/m)^{1/6} v_b$$

$v_b \leq$ (escape velocity from a star) $\sim 600\text{km/s} \Rightarrow v \leq 6000\text{km/s}$

A restricted parabolic three-body problem

Sari, SK & Rossi 2010; SK, Rossi & Sari2012

Binary on parabolic orbit around massive BH

For a large mass ratio: BH / binary $\gg 1$, we can formulate the problem as the motion of a single particle under influence of time - dependent forces.

- Energy gain/loss can be simply rescaled in terms of masses and binary separation

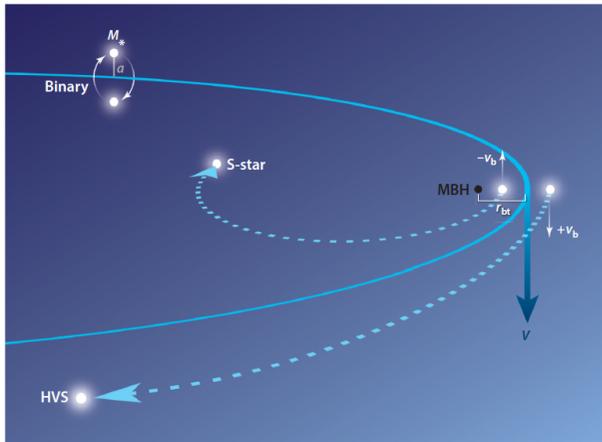
$$\Delta E_1 = -\Delta E_2 = \frac{Gm_1m_2}{a} \left(\frac{M}{m} \right)^{1/3} f(D, \vec{\ell}, \phi)$$

where $D = R_p/R_t$, $\vec{\ell}$: binary orientation, ϕ : binary phase, f is about unit.

$D = \beta^{-1}$: penetration factor

- Analytic solutions for deep penetration $D \ll 1$

- Even for an unequal-mass binary, each member is ejected in exactly 50% of the cases.



Parabolic orbits : zero energy

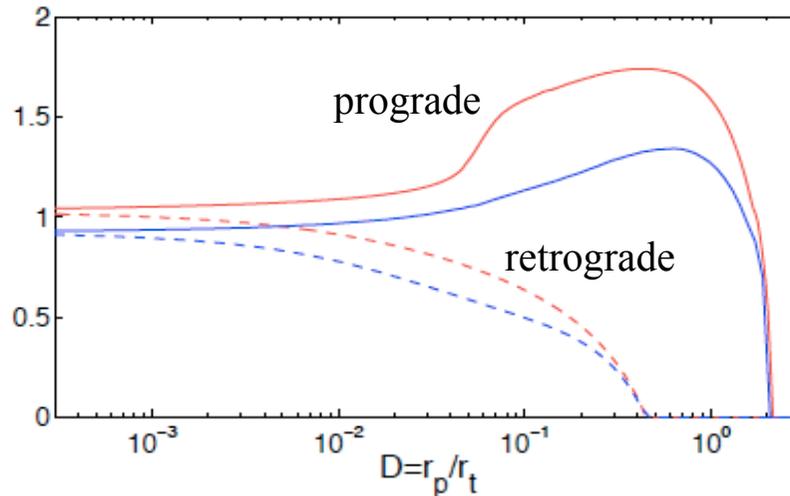
$$\Delta E_1 = -\Delta E_2$$

$$\Delta E_1(\phi + \pi) = -\Delta E_1(\phi)$$

- The radius of gravitational influence of the BH \sim pc $\sim 10^5 R_{\text{tidal}}$ for a \sim several solar radii.
 - equal ejection chance up to $m_1/m_2 \sim 10^2$.

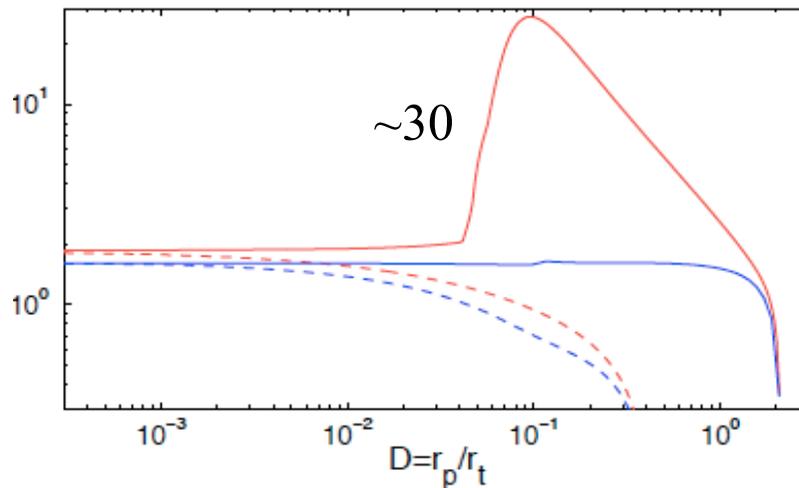
$$\Delta E = \frac{Gm_1m_2}{a} \left(\frac{M}{m} \right)^{1/3} f(D, \vec{\ell}, \phi)$$

$\langle |f| \rangle$



When binaries penetrate the tidal radius deeply, they experience very strong tide. However, ejection energy (ejection velocity) does not become large.

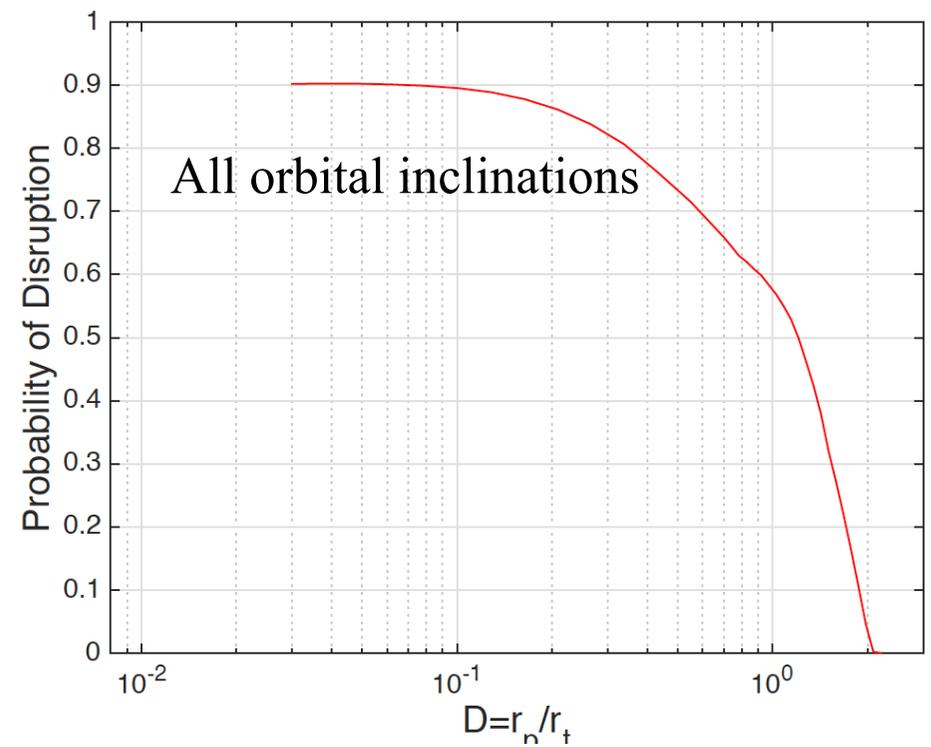
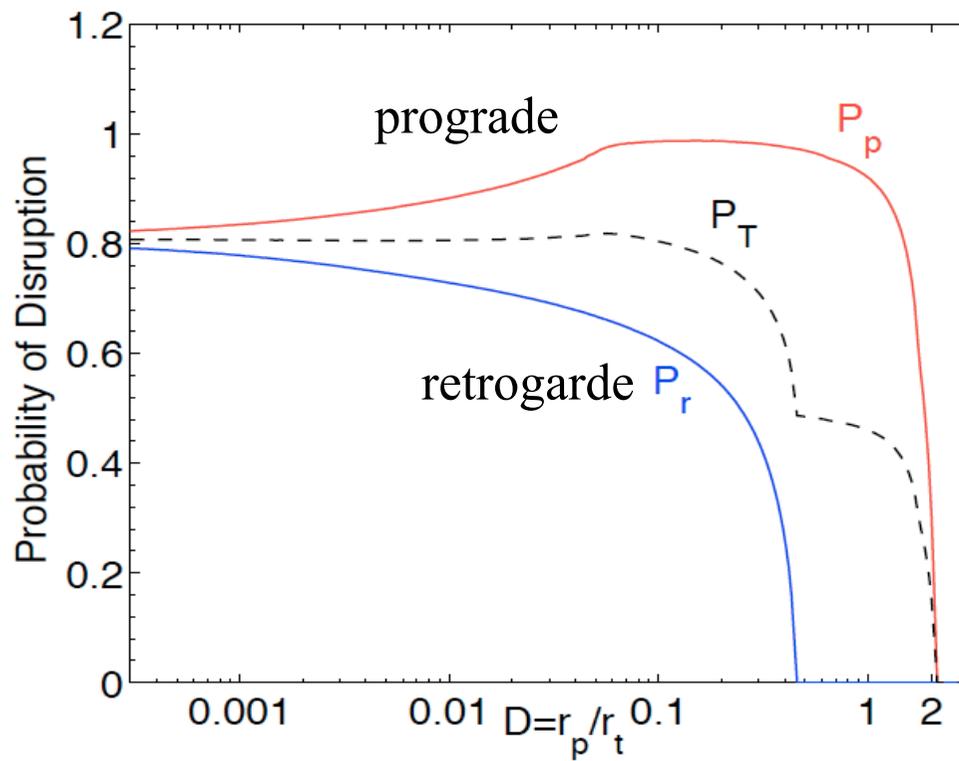
f_{\max}

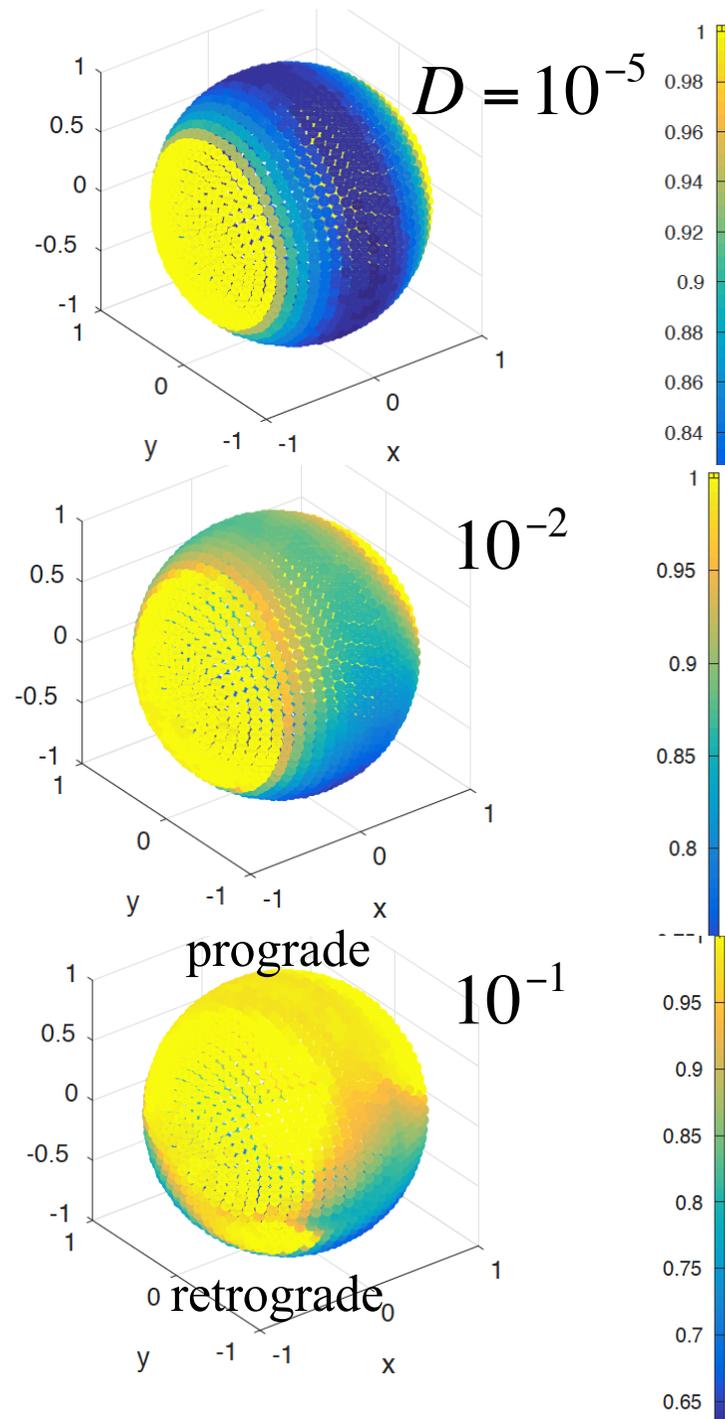
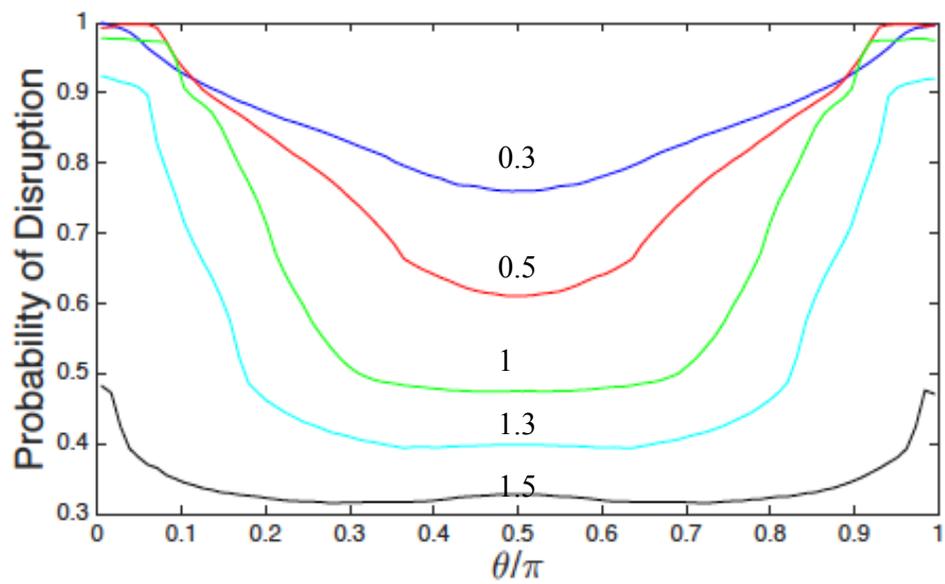
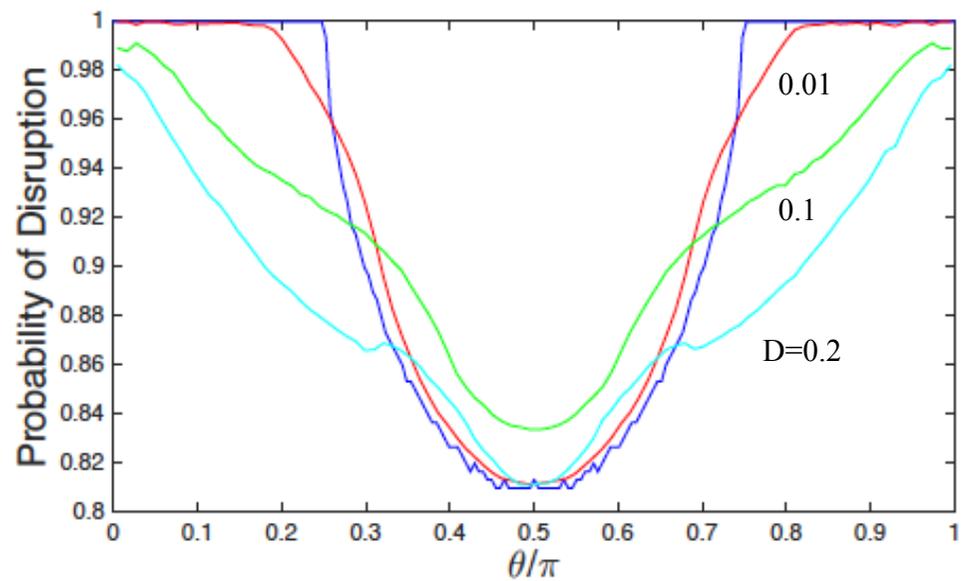


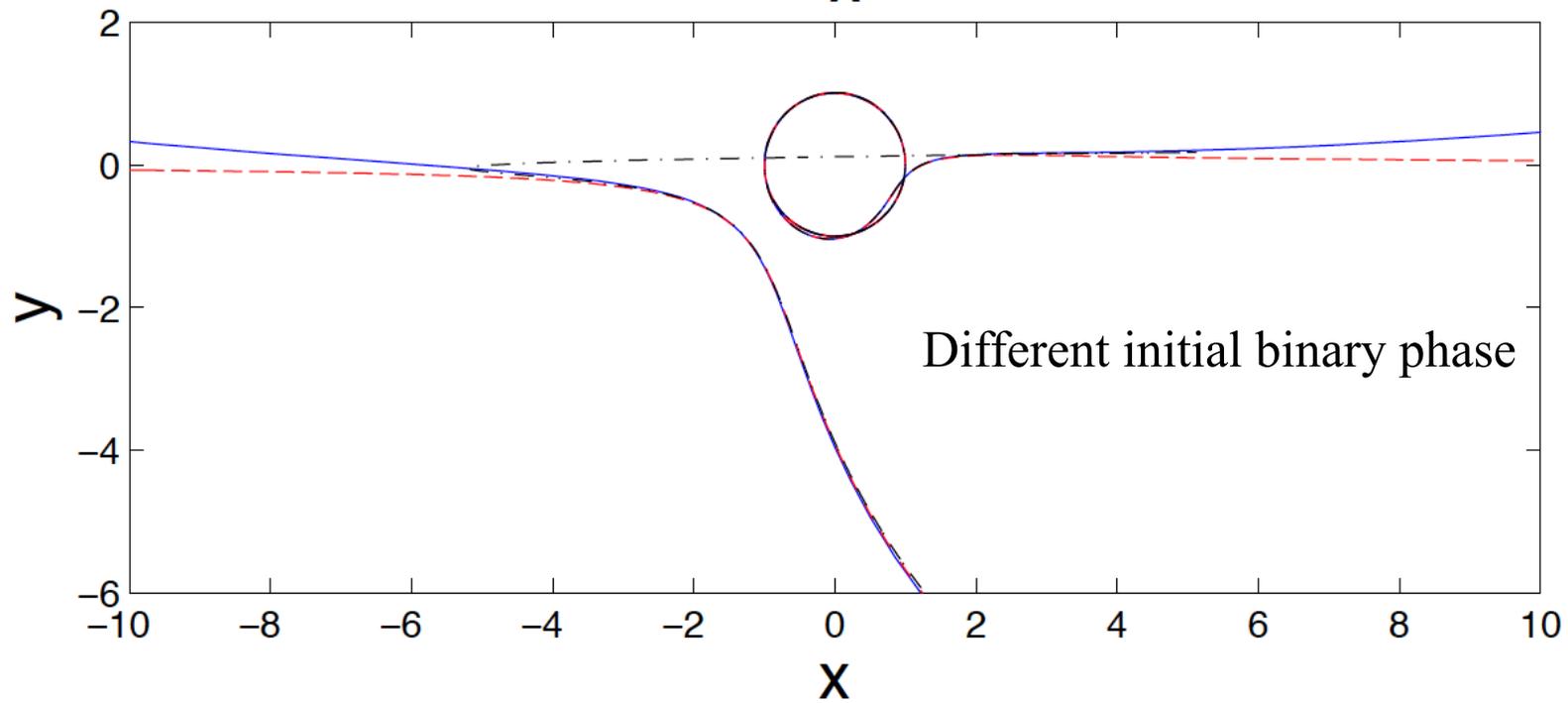
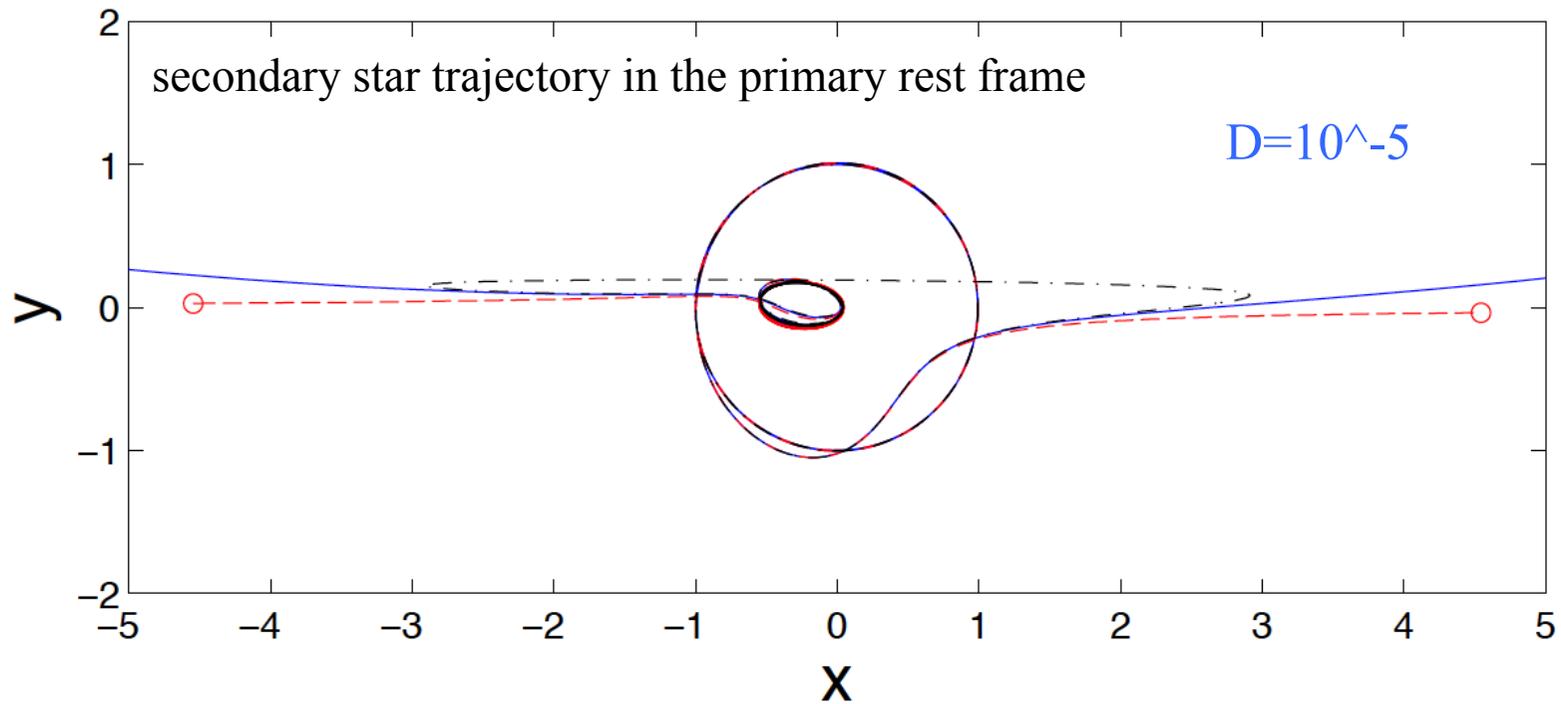
In TDE study, the orbital energy spread of tidal debris is often estimated as $\Delta \varepsilon \sim \frac{GMR_*}{R_p^3}$

but it should be $\sim \frac{GMR_*}{R_t^3}$

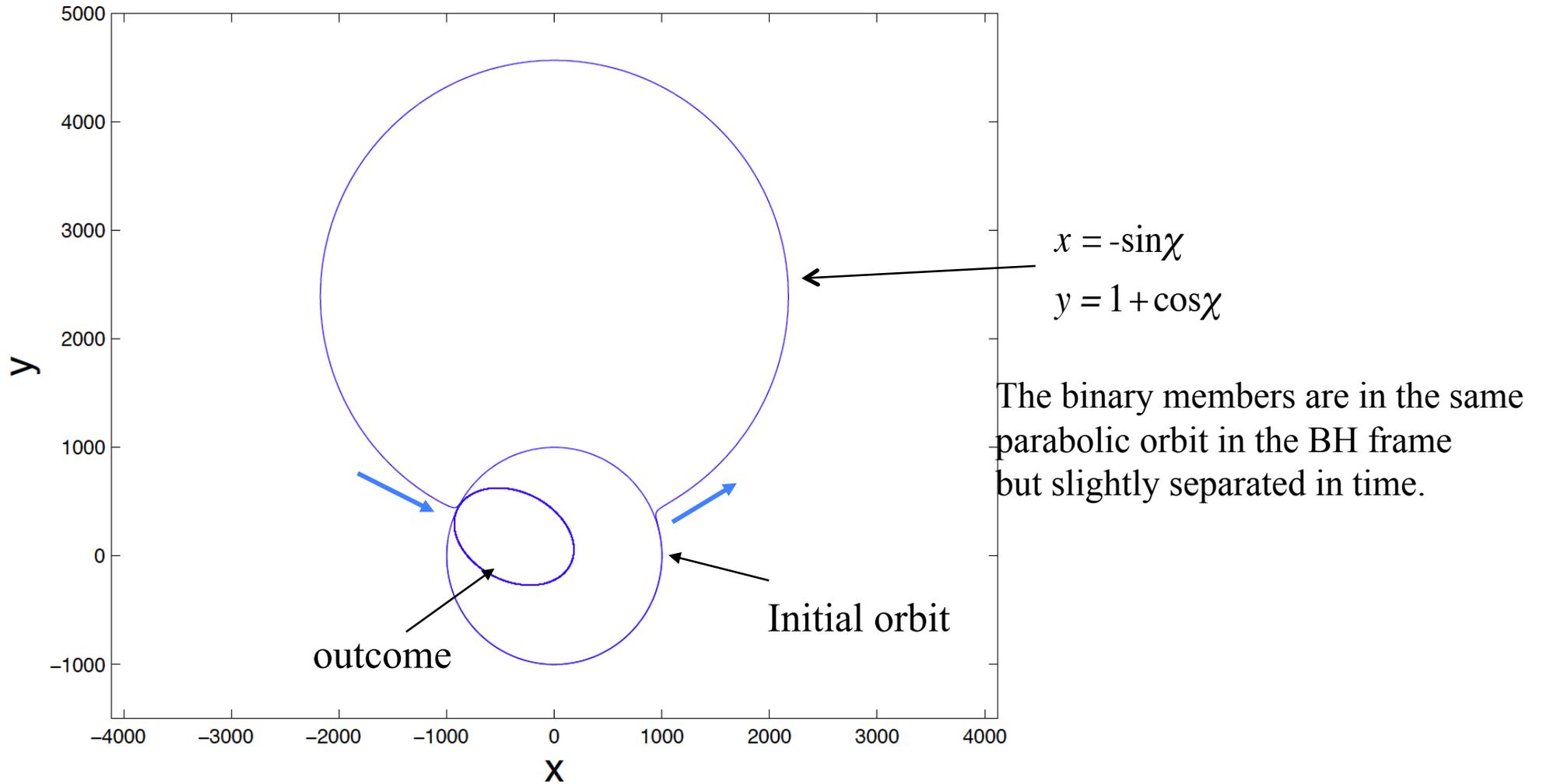
- Even if binaries deeply penetrate the tidal radius, they survive in 10% of the cases, and they become compact binaries.



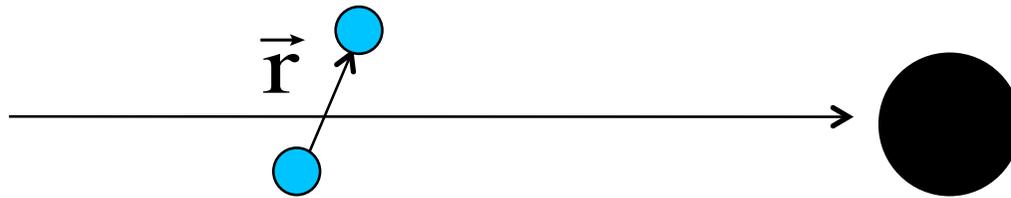




In deep encounters, one of the free solutions dominates around the periapsis passage. Binary dissolves, but after the periapsis passage, they come close to each other.



Radial case



$$\ddot{x} = \frac{4}{9t^2}x - \frac{x}{(x^2 + y^2 + z^2)^{3/2}},$$

$$\ddot{y} = -\frac{2}{9t^2}y - \frac{y}{(x^2 + y^2 + z^2)^{3/2}},$$

$$\ddot{z} = -\frac{2}{9t^2}z - \frac{z}{(x^2 + y^2 + z^2)^{3/2}}.$$

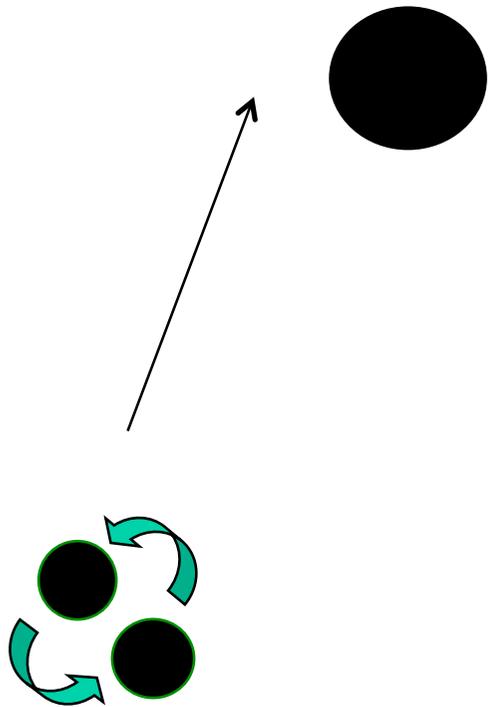
$$x(t) = A_x |t|^{-1/3} + B_x |t|^{4/3},$$

$$y(t) = A_y |t|^{1/3} + B_y |t|^{2/3},$$

$$z(t) = A_z |t|^{1/3} + B_z |t|^{2/3}.$$

Ax solution describes two particles that have the same trajectory, but are slightly separated in time.

Tidal disruption of BH-BH binaries by massive BH



Hypervelocity BHs

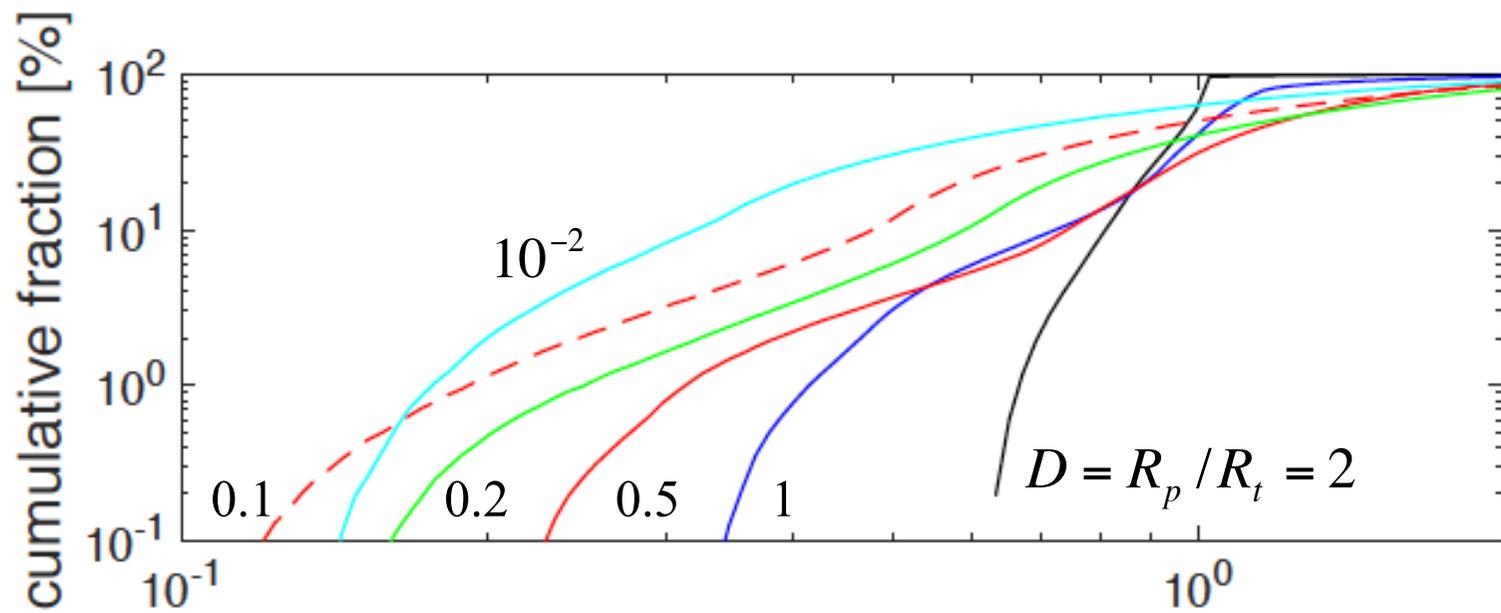
Bound ones: EMRI GW

Survivors

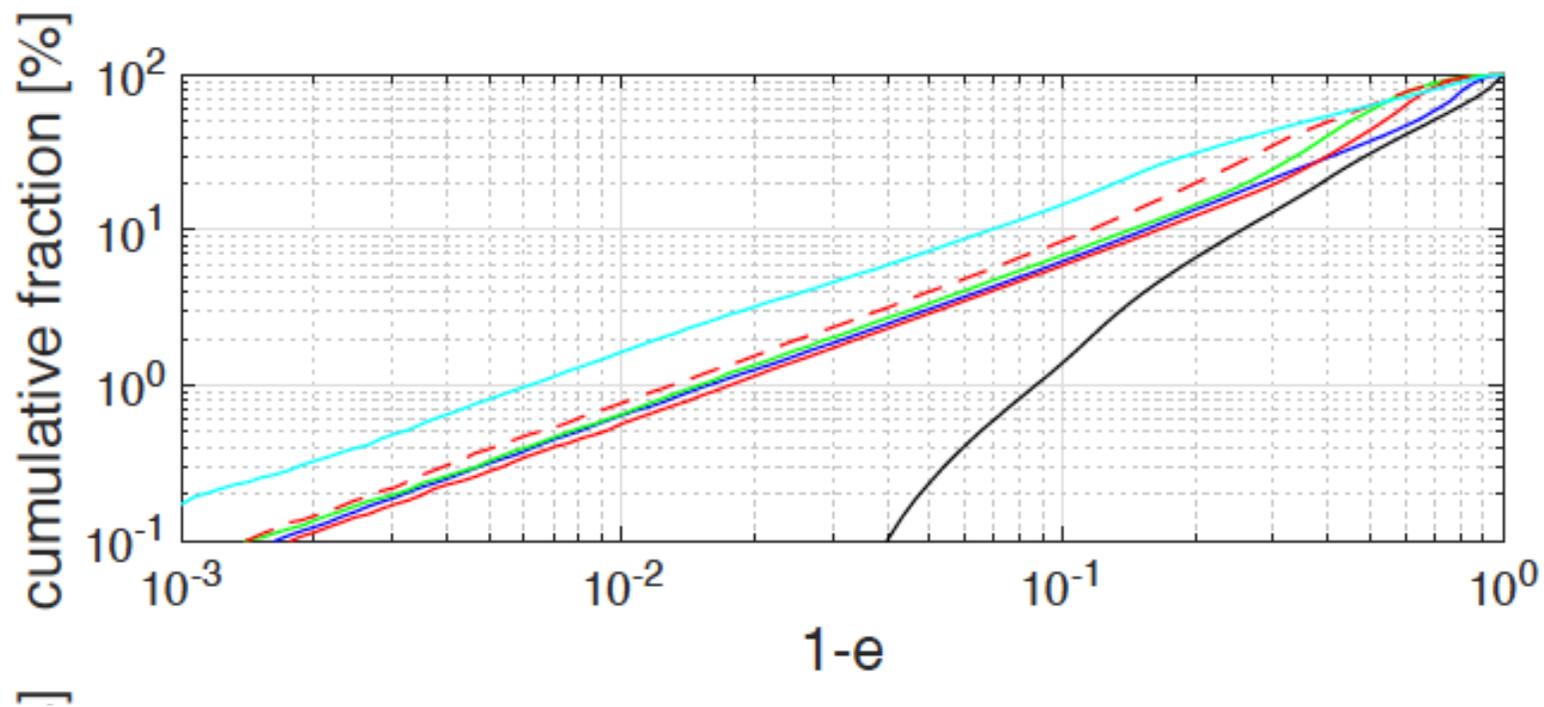
Eccentric, Compact BH binaries

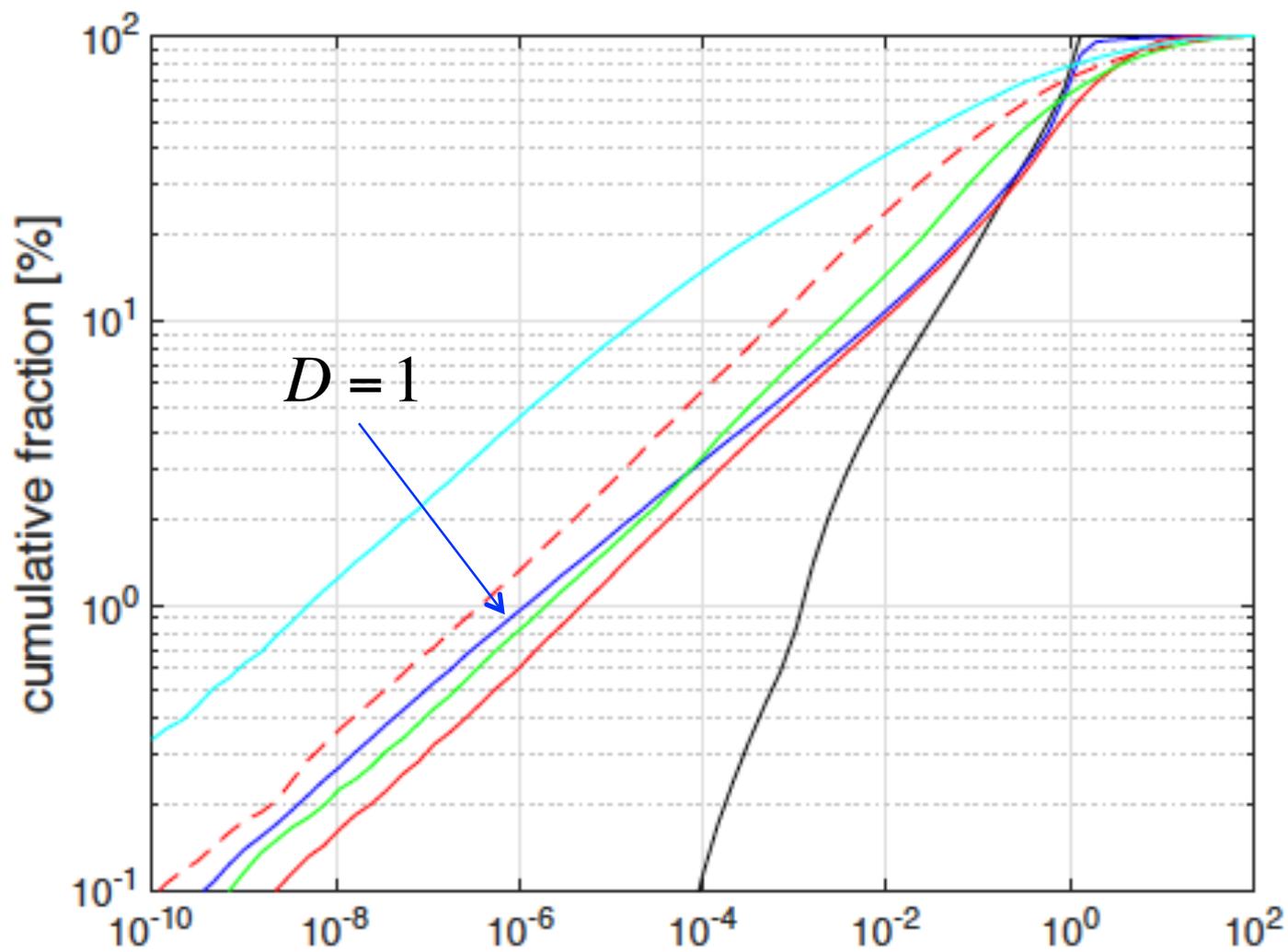
Gravitational wave merger time scale (Peters 1964)

$$T_{GW} = \frac{3}{85} \left(\frac{c^5 a_0^4}{G^3 m_1 m_2 m} \right) (1 - e_0^2)^{7/2}$$
$$= 1.2 \times 10^{14} \text{ yrs} \left(\frac{m}{20 M_{sun}} \right)^{-3} \left(\frac{a_0}{1 \text{ AU}} \right)^4 (1 - e_0^2)^{7/2}$$



a





$$\frac{t_{gw,out}}{t_{gw,in}} = \left(\frac{a_{out}}{a_{in}} \right)^4 (1 - e_{out}^2)^{7/2}$$

Summary

- HVSs: tidal disruption of binaries by a massive BH
- Restricted three-body formula ($M \gg m$)
 - Equal ejection chance even for unequal-mass binaries
 - Deep penetration cases do not provide highest ejection velocity. $D \sim 0.1$
 - Even very deep penetration cases: 10% survive
 - Survivors: smaller semi-major axis, significant eccentricity
- Tidal disruption of BH-BH binaries
 - survivors have much shorter GW merger times.
 - Kozai mechanism: Antonini and Perets 2012; VanLandingham et al. 2016
- GAIA Data Release 1 on 14 Sep 2016.
 - Positions, **proper motions** and parallaxes for 1 billion stars; nearby (10kpc) HVSs
 - Accurate proper motion: key for constraining the Galactic potential.